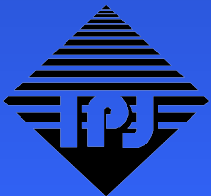


Single-particle effects in the properties of heavy and superheavy nuclei

Aleksander Parkhomenko



Plan

- Introduction
- Method of calculations
- Phenomenological formula for α -decay half-lives
- Results
 - a) One-quasiparticle excitations
 - b) Systematic
 - c) α -decay chains
- Conclusions

MENDELEYEV PERIODIC TABLE OF THE ELEMENTS

1 H																	2 He
3 Li	4 Be											5 B	6 C	7 N	8 O	9 F	10 Ne
11 Na	12 Mg											13 Al	14 Si	15 P	16 S	17 Cl	18 Ar
19 K	20 Ca	21 Sc	22 Ti	23 V	24 Cr	25 Mn	26 Fe	27 Co	28 Ni	29 Cu	30 Zn	31 Ga	32 Ge	33 As	34 Se	35 Br	36 Kr
37 Rb	38 Sr	39 Y	40 Zr	41 Nb	42 Mo	43 Tc	44 Ru	45 Rh	46 Pd	47 Ag	48 Cd	49 In	50 Sn	51 Sb	52 Te	53 I	54 Xe
55 Cs	56 Ba	57 La*	72 Hf	73 Ta	74 Re	75 Os	76 Ir	78 Pt	79 Au	80 Hg	81 Tl	82 Pb	83 Bi	84 Po	85 At	86 Rn	
87 Fr	88 Ra	89 Ac+	105 Db	106 Bh	108 Hs	109 Mt	110 Ds	111 Rg	112 Cn	113	114	115	116				118

105 Db

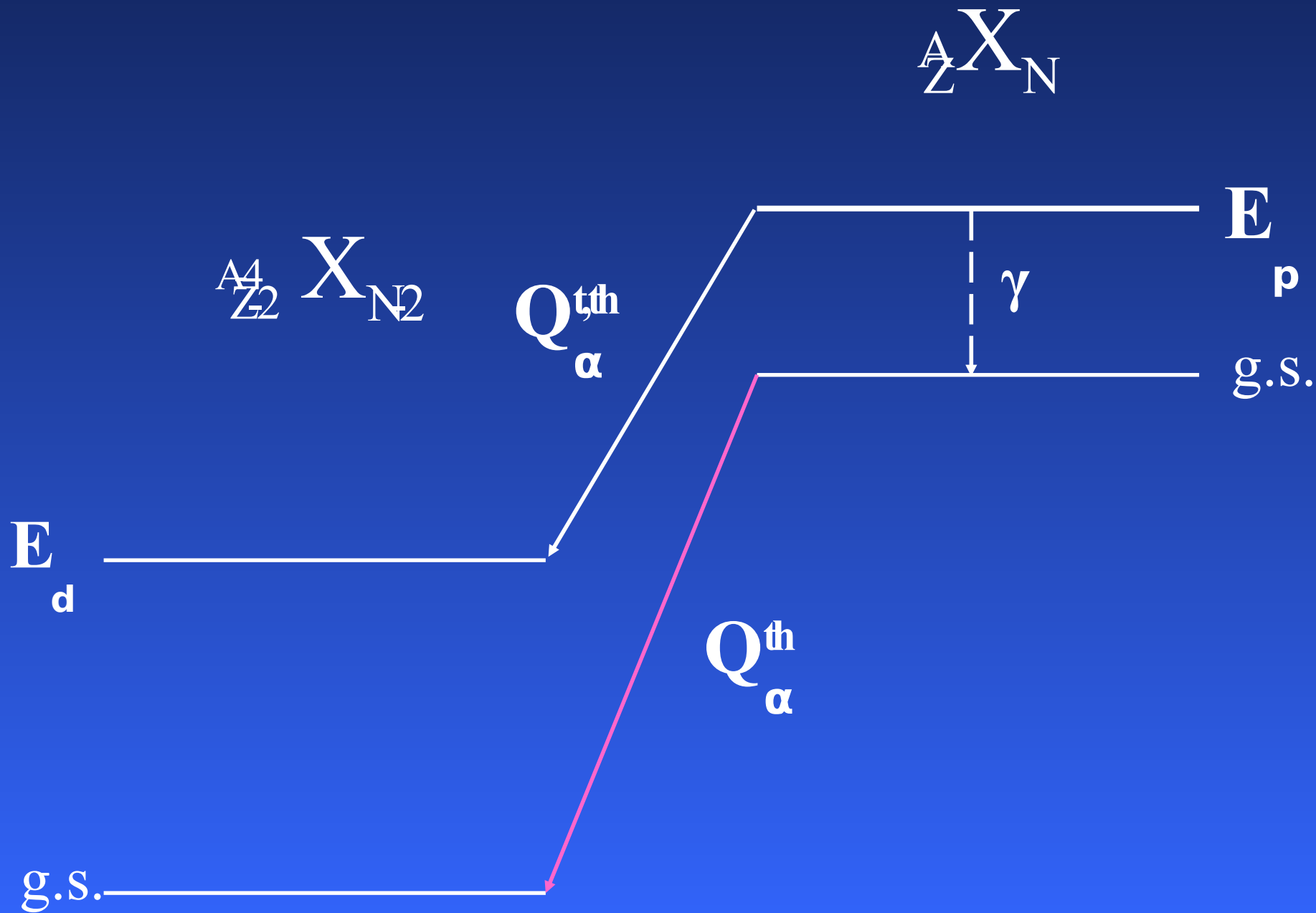
Transactinides

+ Actinides

90 Th	91 Pa	92 U	93 Np	94 Pu	95 Am	96 Cm	97 Bk	98 Cf	99 Es	100 Fm	101 Md	102 No	103 Lr
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*** Lanthanides**

58 Ce	59 Pr	60 Nd	61 Pm	62 Sm	63 Eu	64 Gd	65 Tb	66 Dy	67 Ho	68 Er	69 Tm	70 Yb	71 Lu
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Method of the calculations

Macroscopic

Microscopic

Yukawa-plus-exponential
model

Strutinski shell correction

BCS pairing
approximation

Woods-Saxon potential

The shape of a nucleus

$$R(\theta, \phi) = R_0 C(\beta_{\lambda\mu}) \left(1 + \sum_{\lambda\mu} \beta_{\lambda\mu} Y_{\lambda\mu}(\theta, \phi) \right)$$

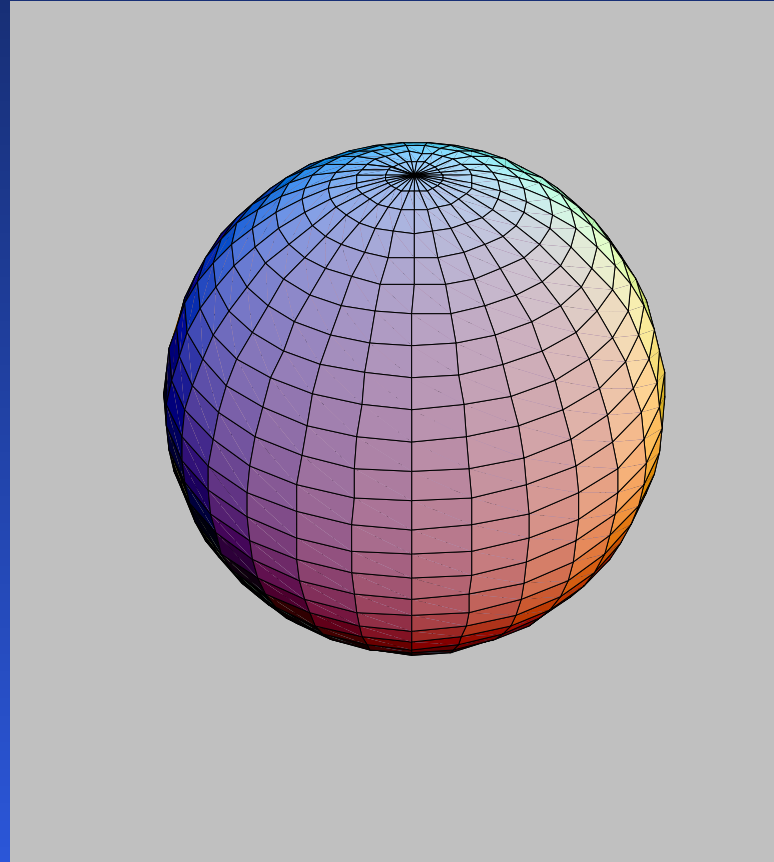
Axial symmetric nuclei:

$$\mu = 0 \Rightarrow$$

$$\beta_{\lambda 0} \equiv \beta_{\lambda}$$

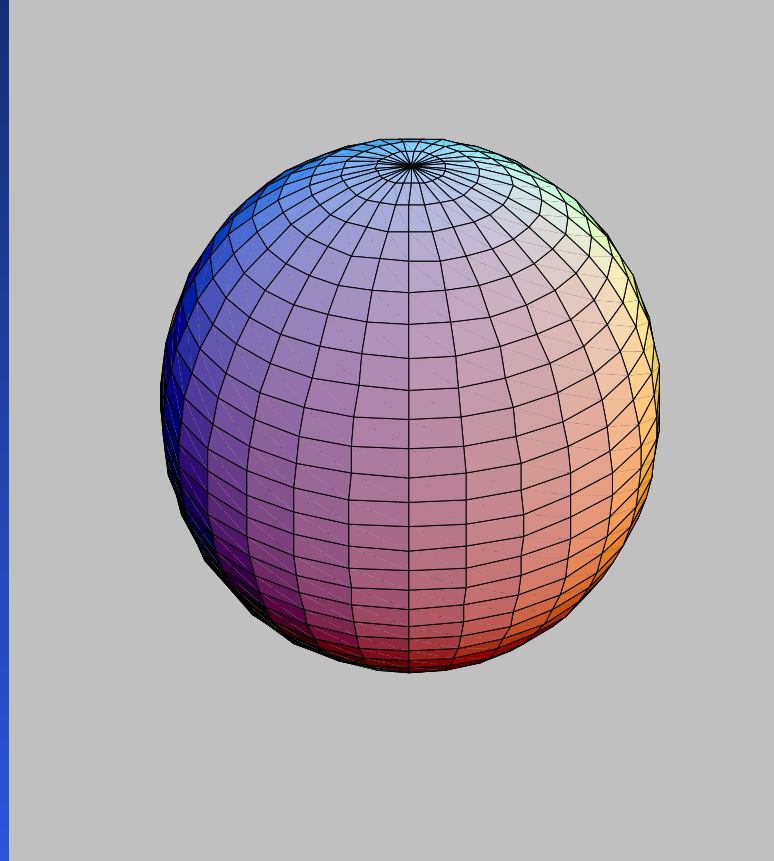
$$R(\theta) = R_0 C(\beta_{\lambda}) \left(1 + \sum_{\lambda=2}^8 \beta_{\lambda} Y_{\lambda 0}(\theta) \right)$$

$$\beta_2 = 0$$



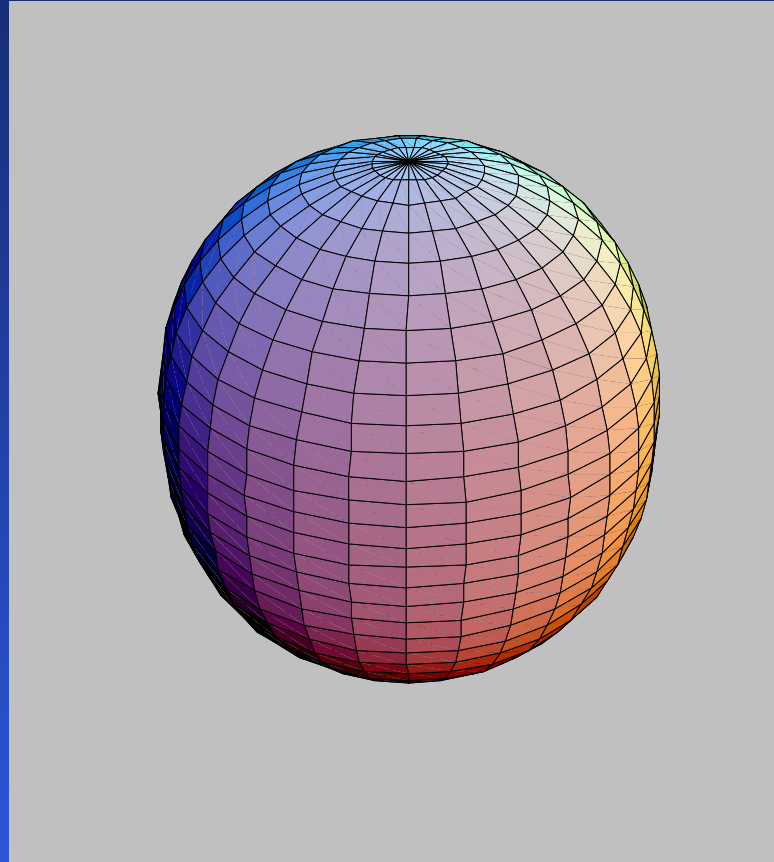
$$\beta_3 = \beta_4 = \beta_5 = \beta_6 = \beta_7 = \beta_8 = 0.$$

$$\beta_2 = 0.2$$



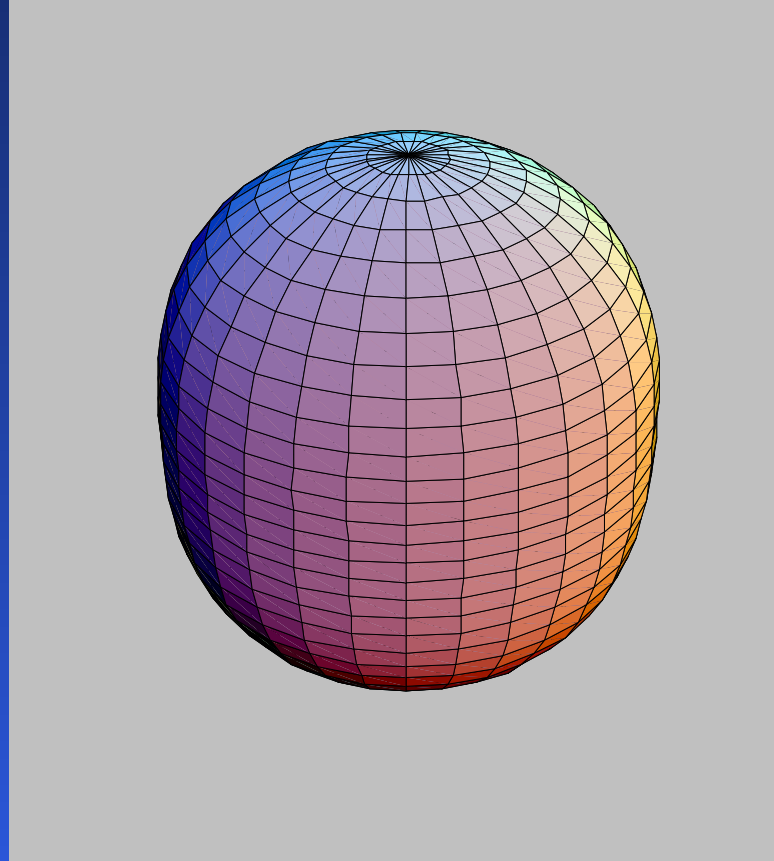
$$\beta_3 = \beta_4 = \beta_5 = \beta_6 = \beta_7 = \beta_8 = 0.$$

$$\beta_2 = 0.3$$



$$\beta_3 = \beta_4 = \beta_5 = \beta_6 = \beta_7 = \beta_8 = 0.$$

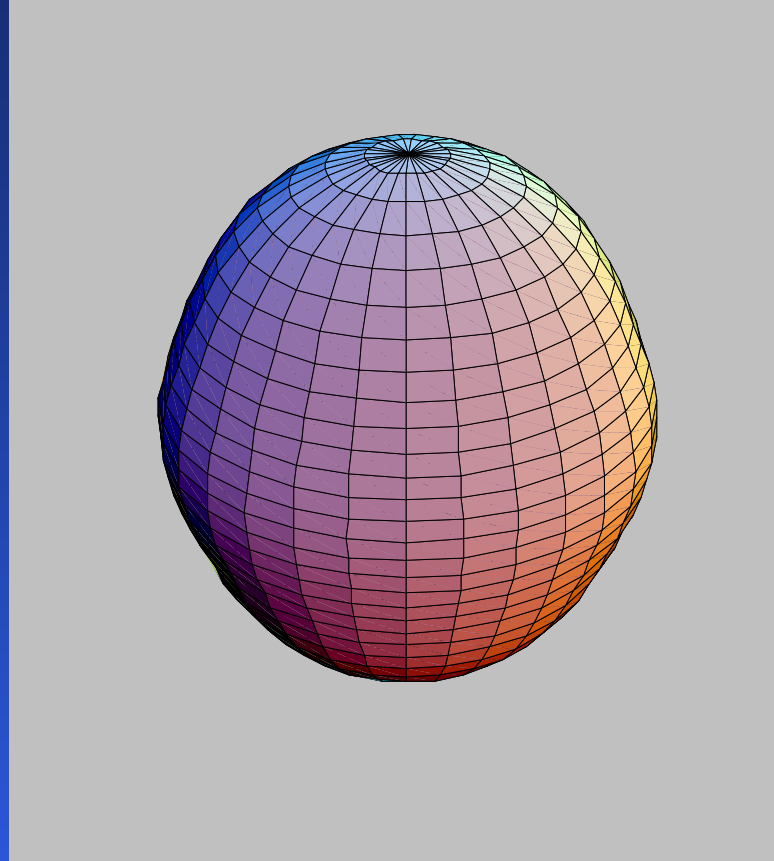
$$\beta_2 = 0.4$$



$$\beta_3 = \beta_4 = \beta_5 = \beta_6 = \beta_7 = \beta_8 = 0.$$

$$\beta_2 = 0.4$$

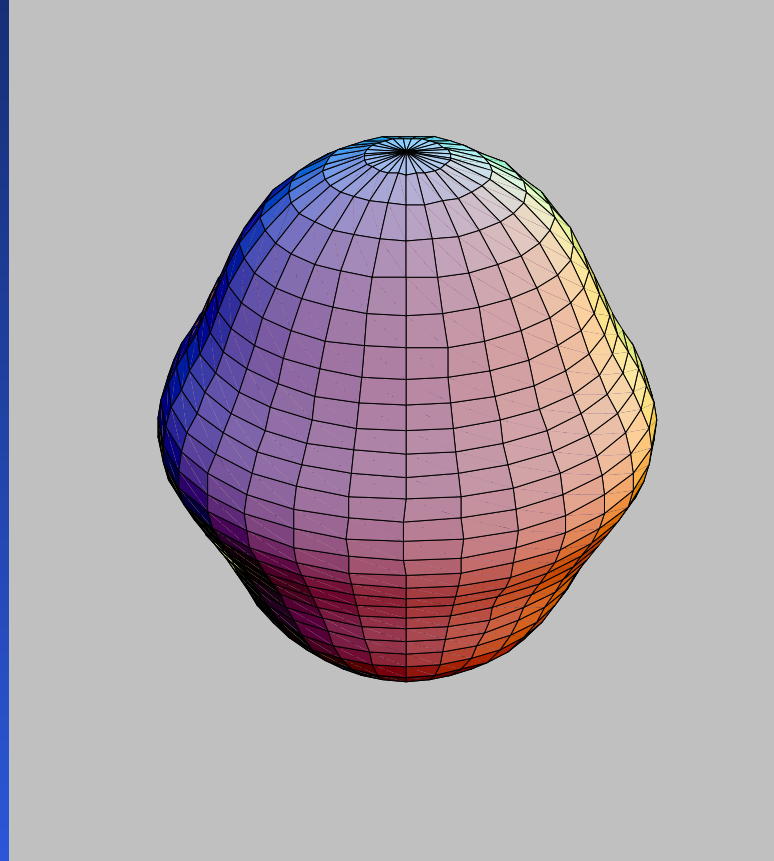
$$\beta_4 = 0.1$$



$$\beta_3 = \beta_5 = \beta_6 = \beta_7 = \beta_8 = 0.$$

$$\beta_2 = 0.4$$

$$\beta_4 = 0.2$$

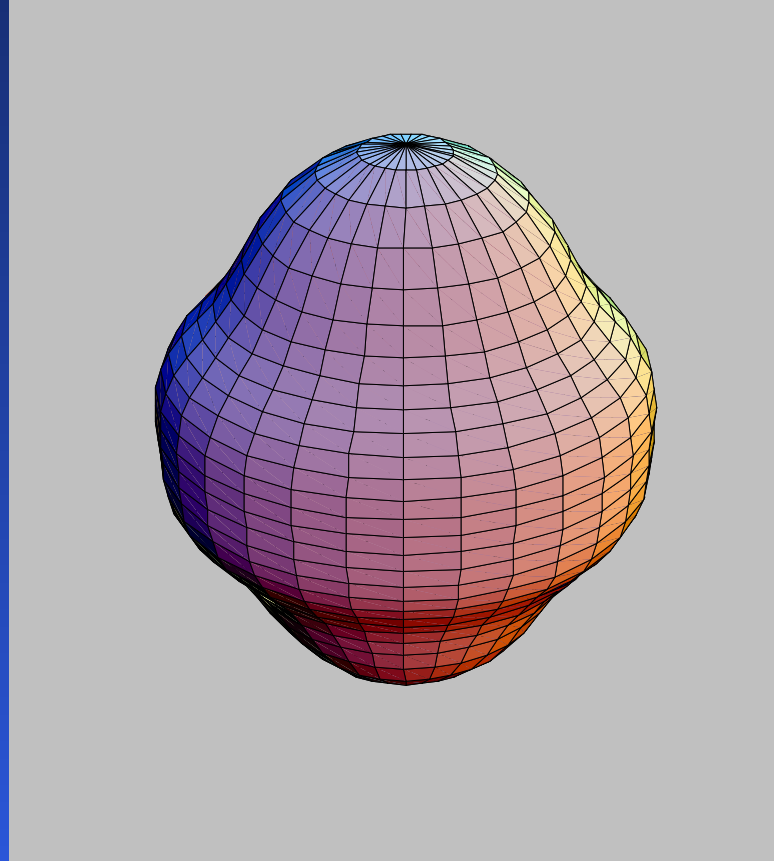


$$\beta_3 = \beta_5 = \beta_6 = \beta_7 = \beta_8 = 0.$$

$$\beta_2 = 0.4$$

$$\beta_4 = 0.2$$

$$\beta_6 = 0.1$$



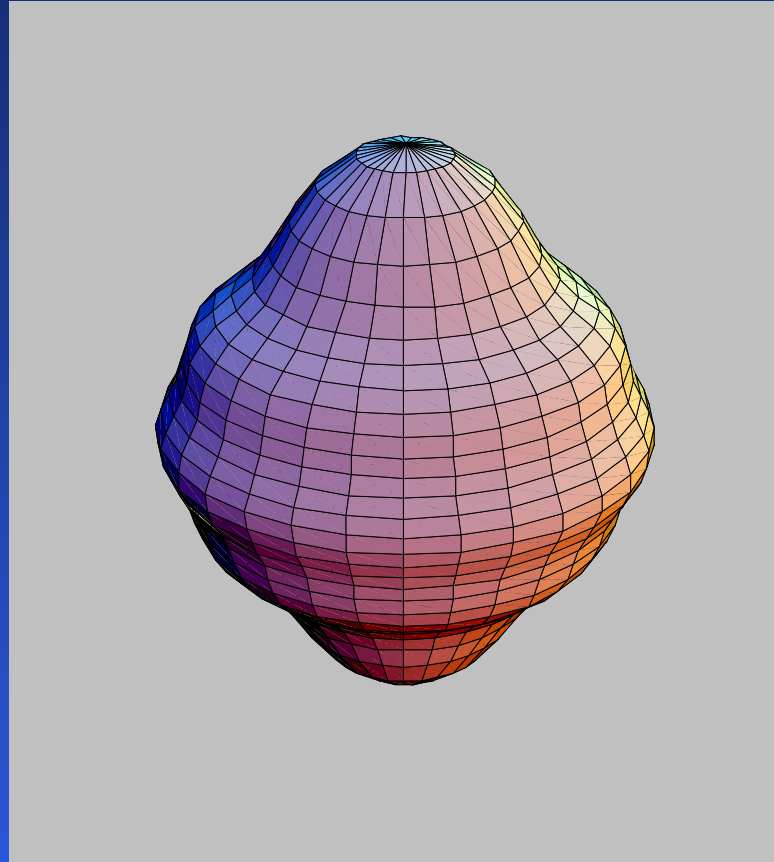
$$\beta_3 = \beta_5 = \beta_7 = \beta_8 = 0$$

$$\beta_2 = 0.4$$

$$\beta_4 = 0.2$$

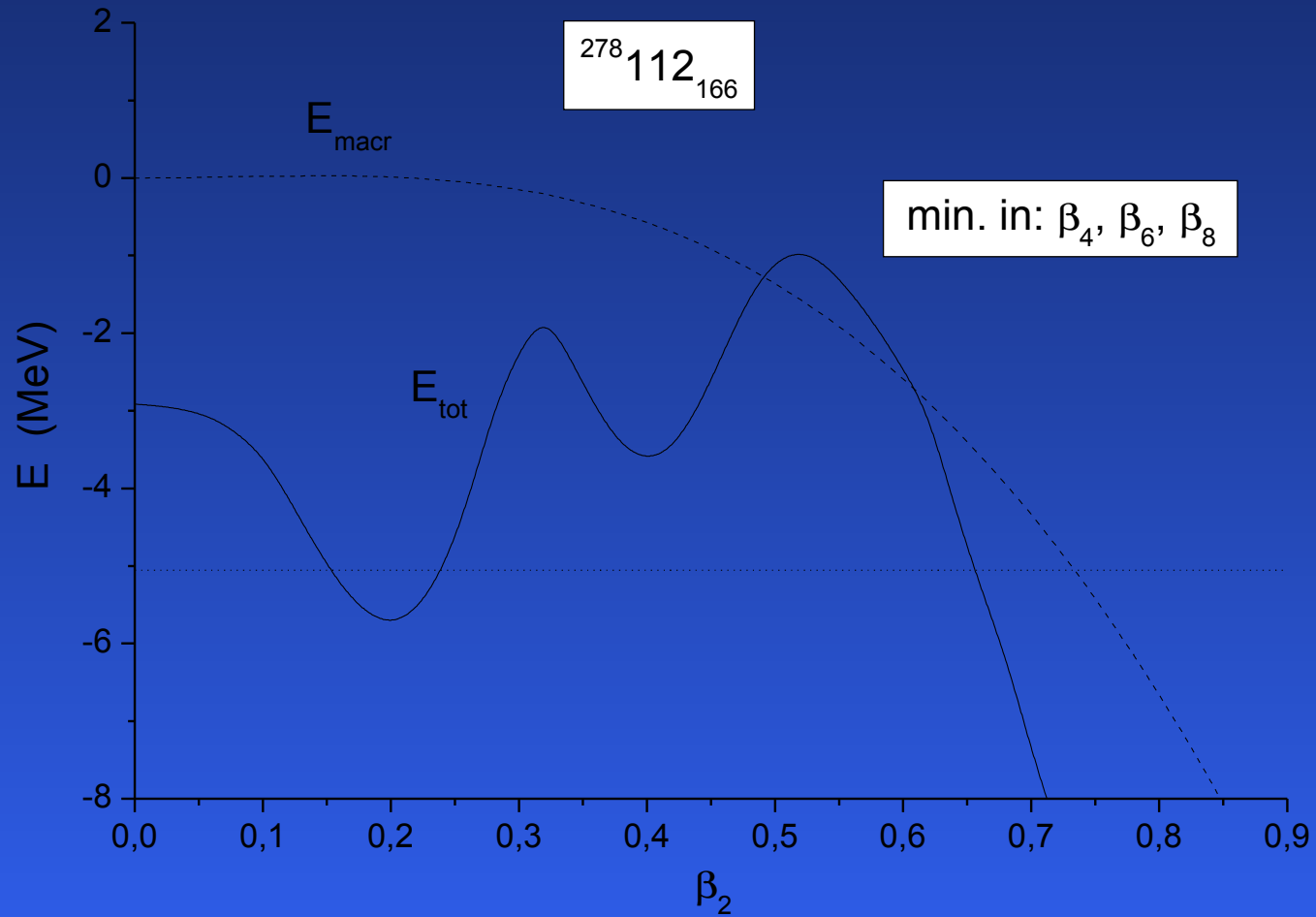
$$\beta_6 = 0.1$$

$$\beta_8 = 0.1$$

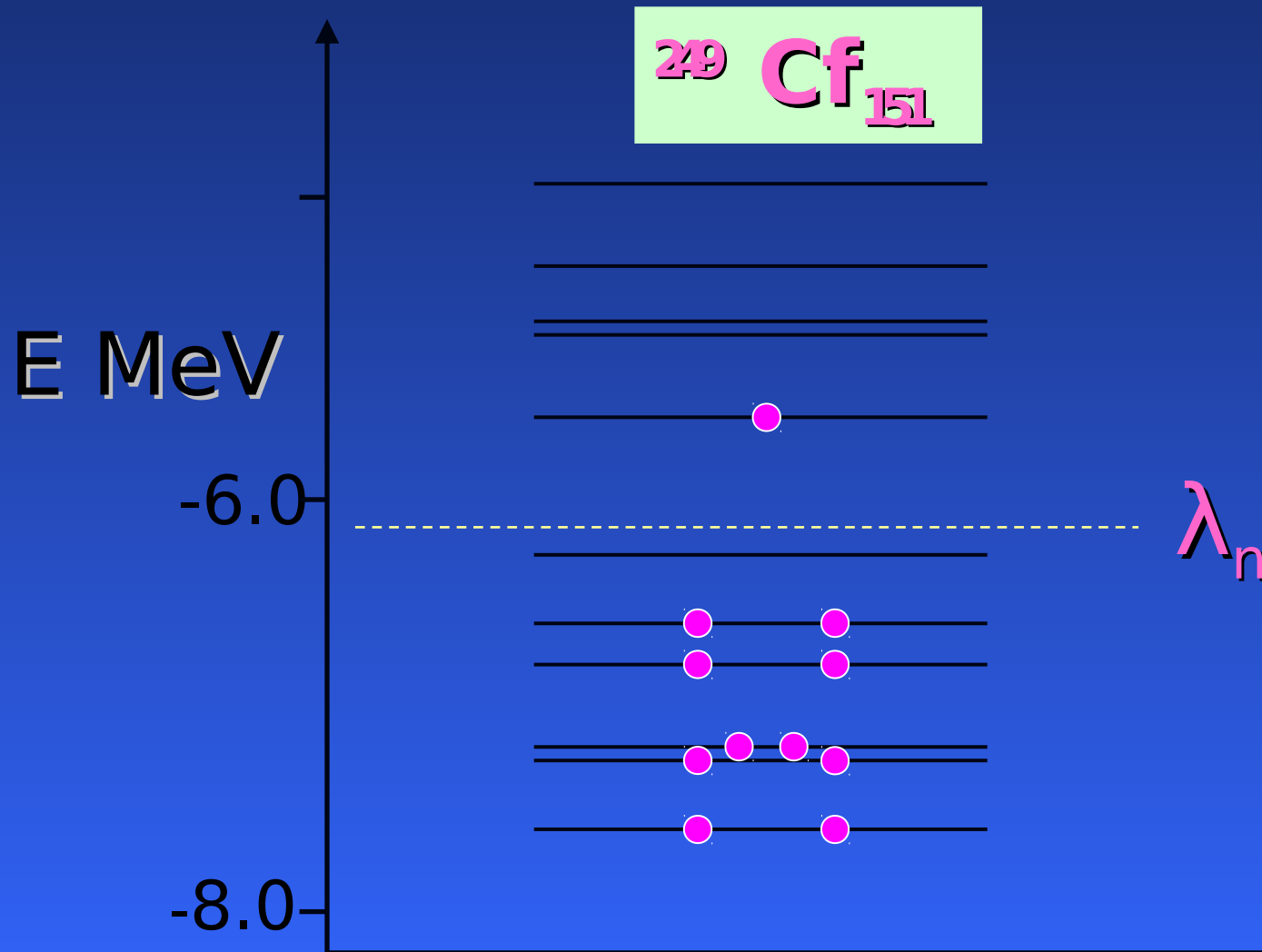


$$\beta_3 = \beta_5 = \beta_7 = 0.$$

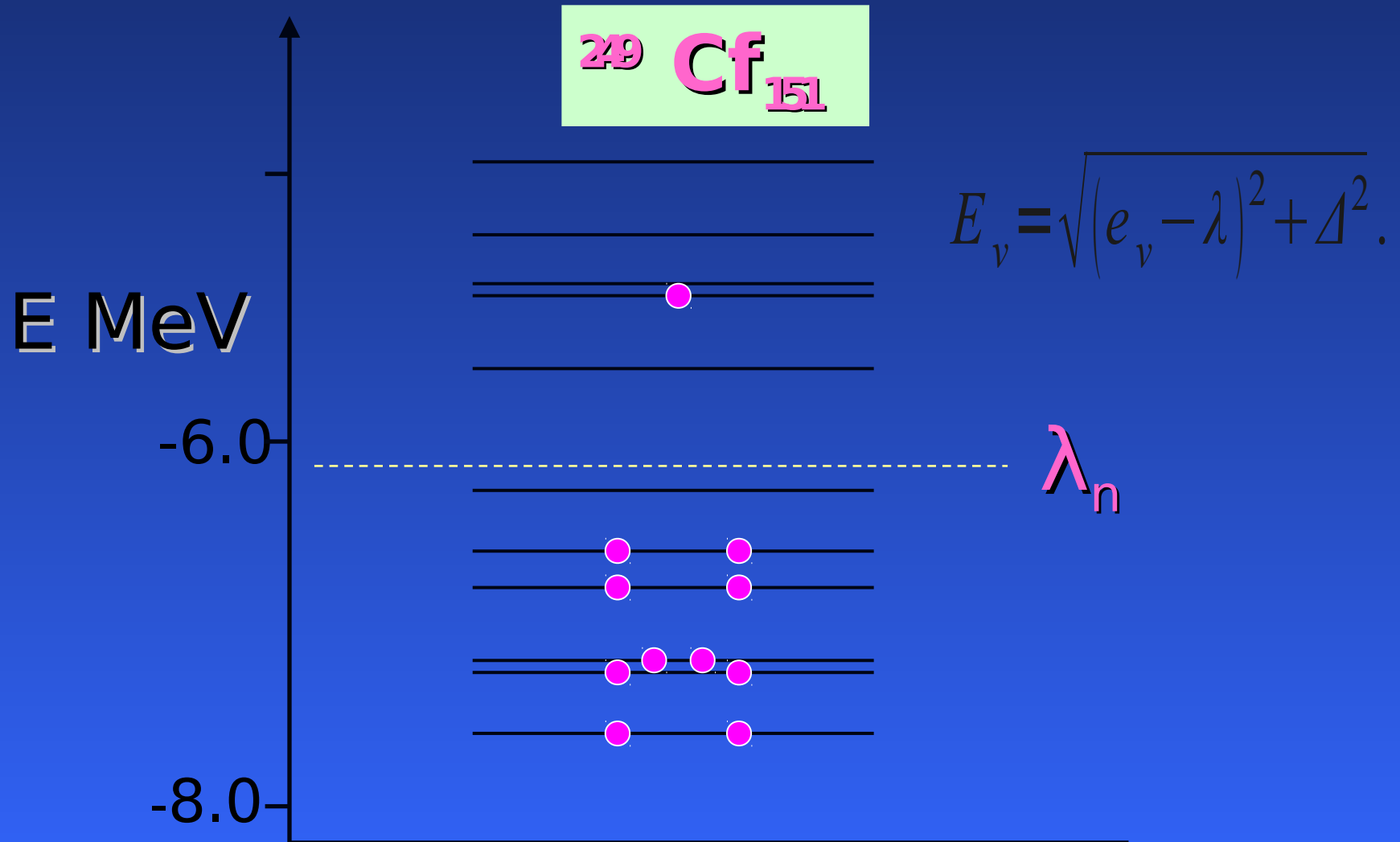
Shell effects



One-quasiparticle excitations



One-quasiparticle excitations



Phenomenological formula for α -decay half-lives

Viola-Seaborg formula :

$$\log_{10} T_{\alpha}(Z, N) = a Z (Q_{\alpha} - E_p)^{-1/2} + b Z + c \quad \text{for e-e}$$

$$\log_{10} T_{\alpha}(Z, N) = a Z (Q_{\alpha} - E_p)^{-1/2} + b Z + c \quad \text{for o-e}$$

$$\log_{10} T_{\alpha}(Z, N) = a Z (Q_{\alpha} - E_n)^{-1/2} + b Z + c \quad \text{for e-o}$$

$$\log_{10} T_{\alpha}(Z, N) = a Z (Q_{\alpha} - E_{mp})^{-1/2} + b Z + c \quad \text{for o-o}$$

$$E_{mp} = E_n + E_p$$

Results

$$\left(\left| \lg T_a^{th}(Q_a^{exp}) - \lg T_a^{exp} \right| \right) \quad \bar{f} = 10^{|\bar{\delta}|}$$

Z = 84 -110 N = 128 -171

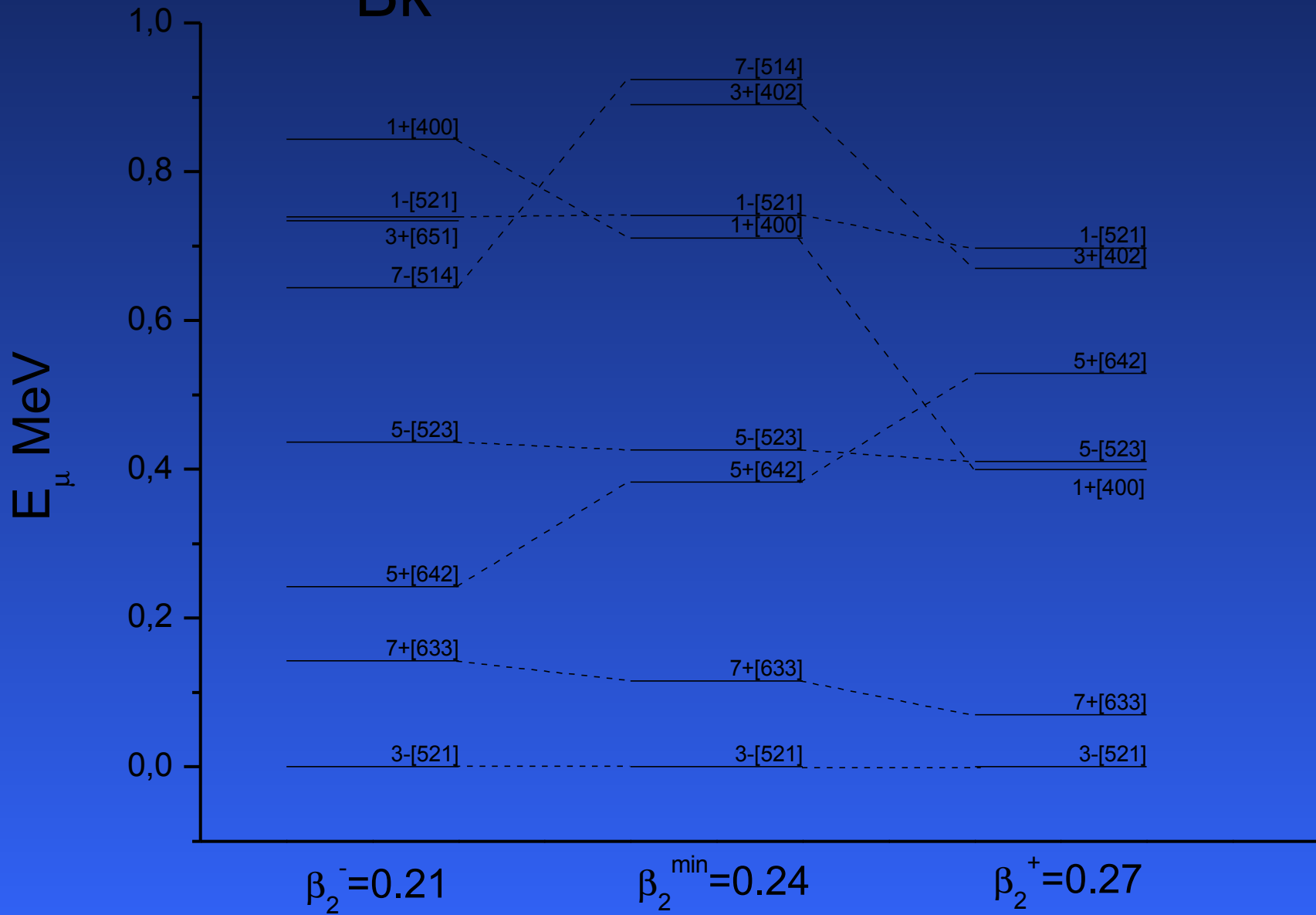
Nuclei	n	$ \bar{\delta} $	rms	\bar{f}	np	\bar{E}
e-e	61	0.13	0.16	1.3	3	0
o-e	45	0.32	0.41	2.1	1	0.11 3
e-o	56	0.50	0.60	3.2	1	0.17 0
o-o	40	0.60	0.72	4.0	0	0.28

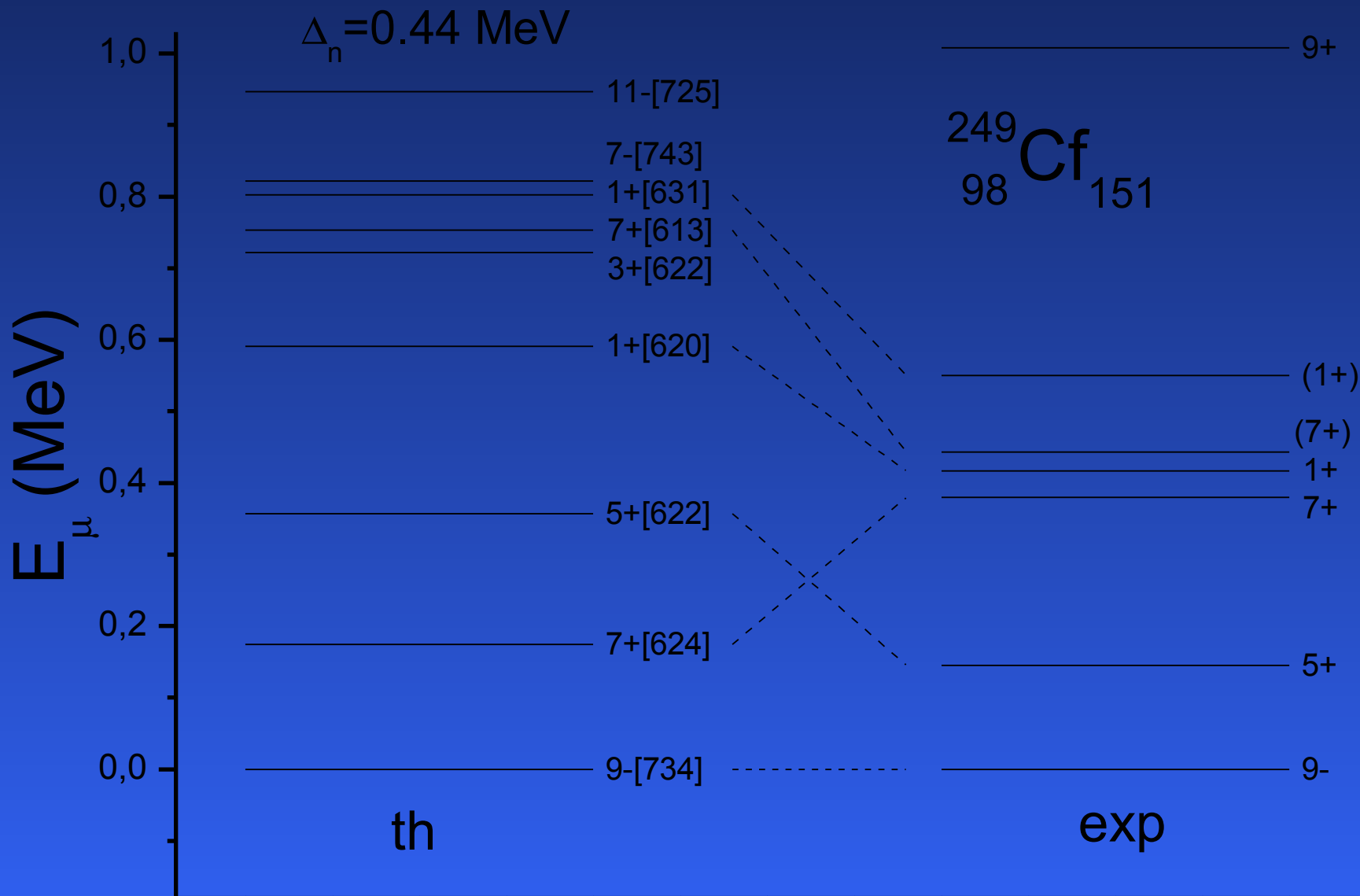
$$\left(\left| \lg T_{\alpha}^{th}(Q_{\alpha}^{th}) - \lg T_{\alpha}^{exp} \right| \right)$$

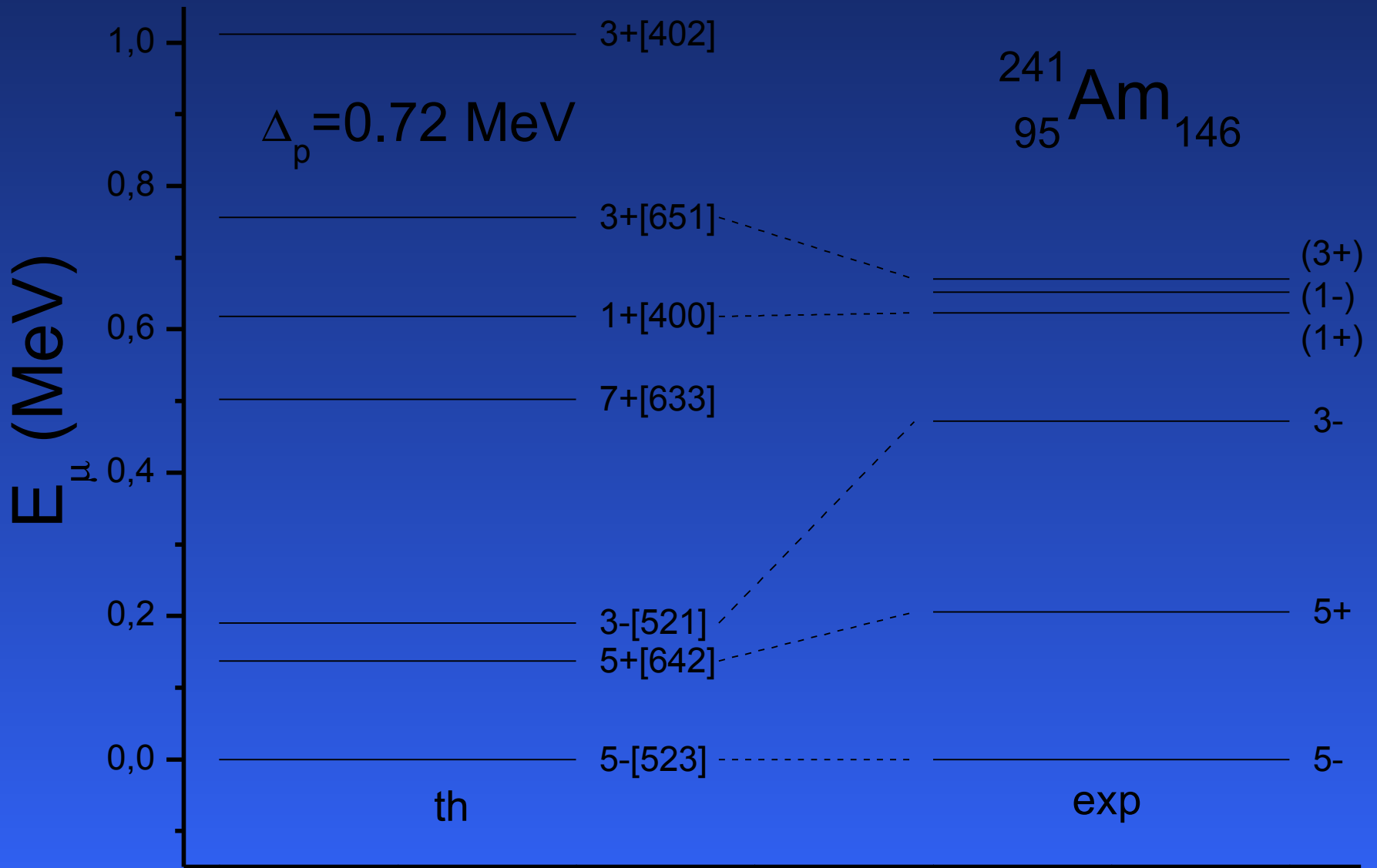
Z = 84 -110 N = 128 -171

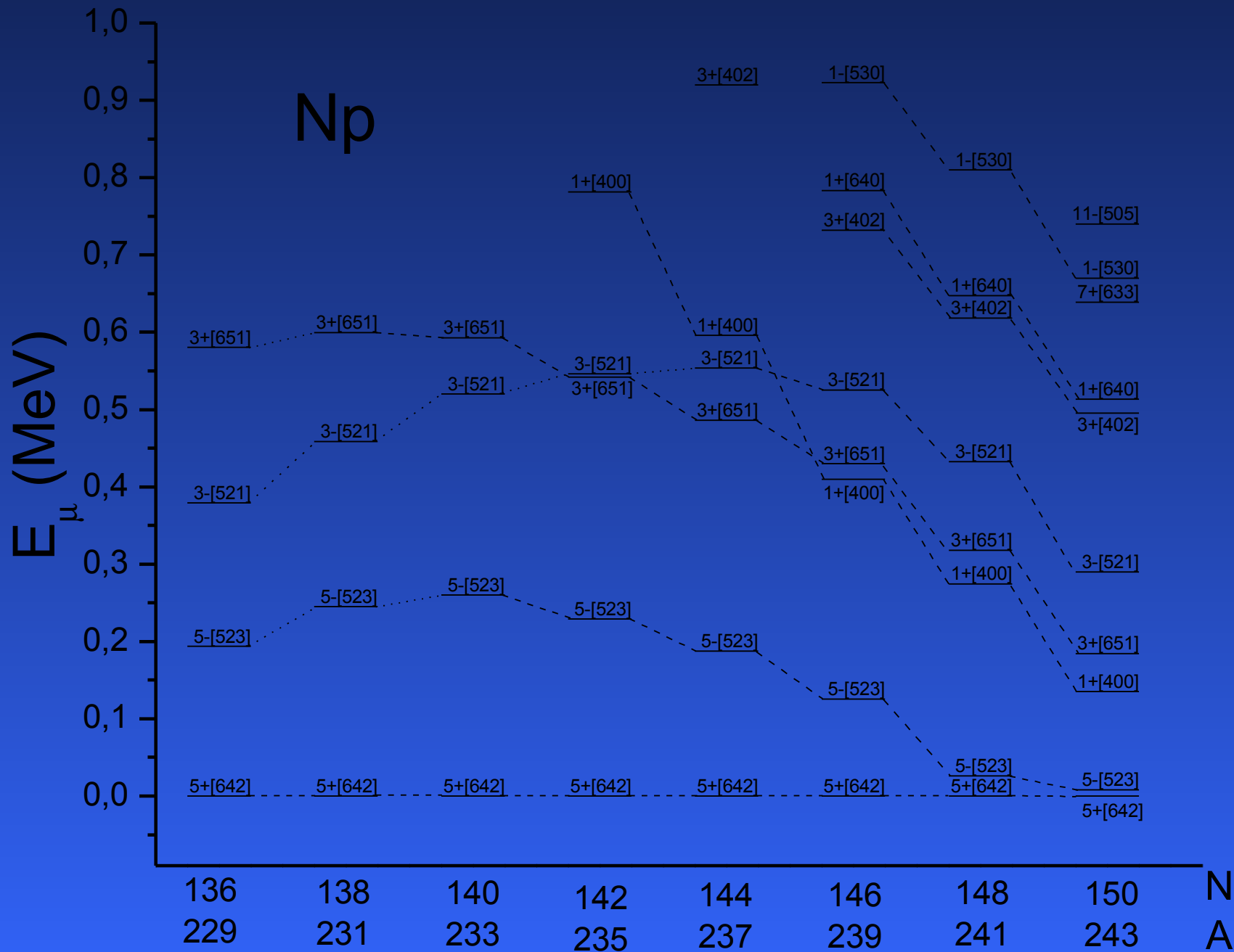
Nuclei	n	$\bar{\delta}$	rms	\bar{f}
e-e	67	0.71	0.93	5.1
o-e	48	1.08	1.44	12.0
e-o	68	1.20	1.58	15.8
o-o	43	1.26	1.57	18.0

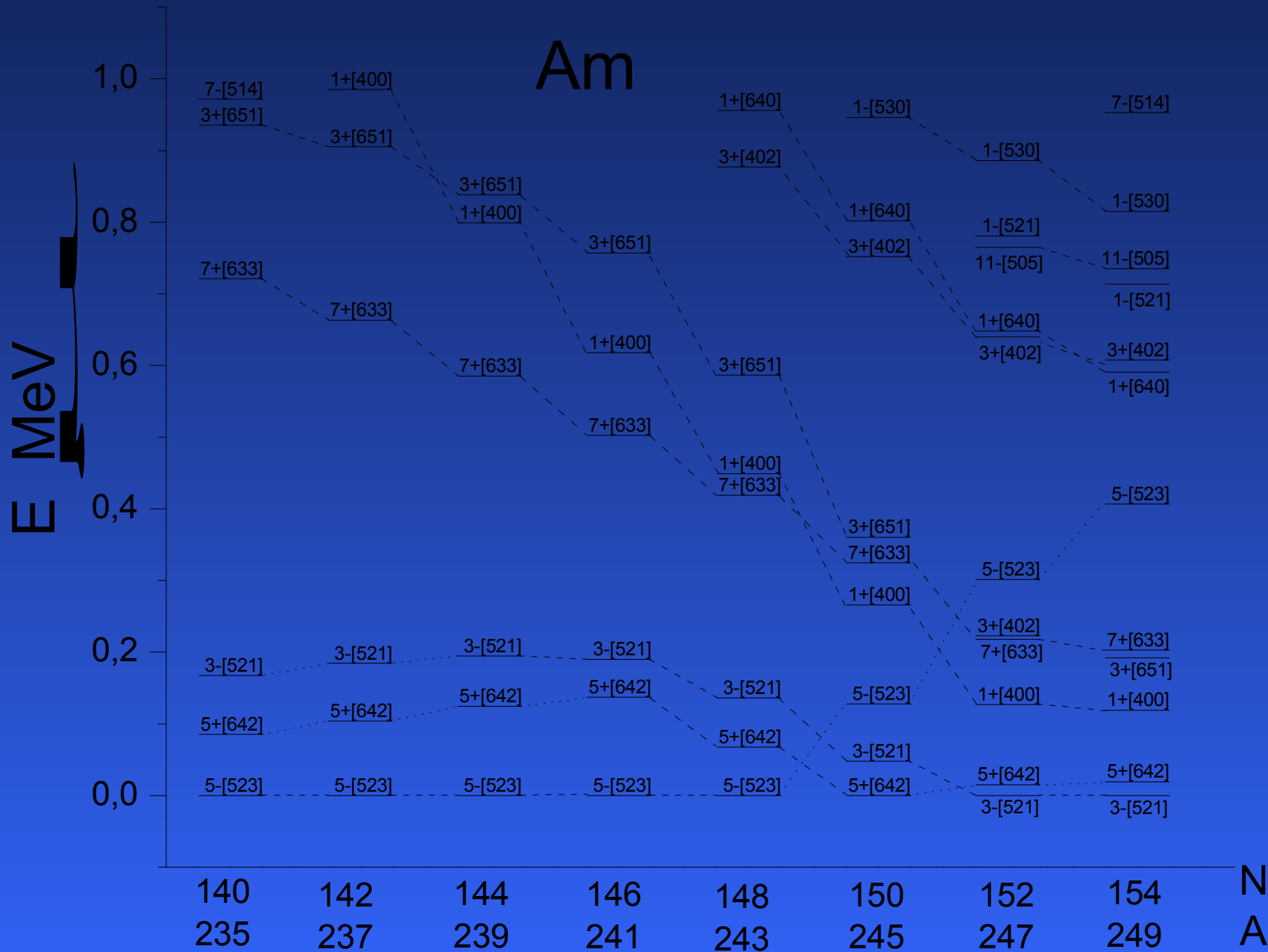
^{245}Bk



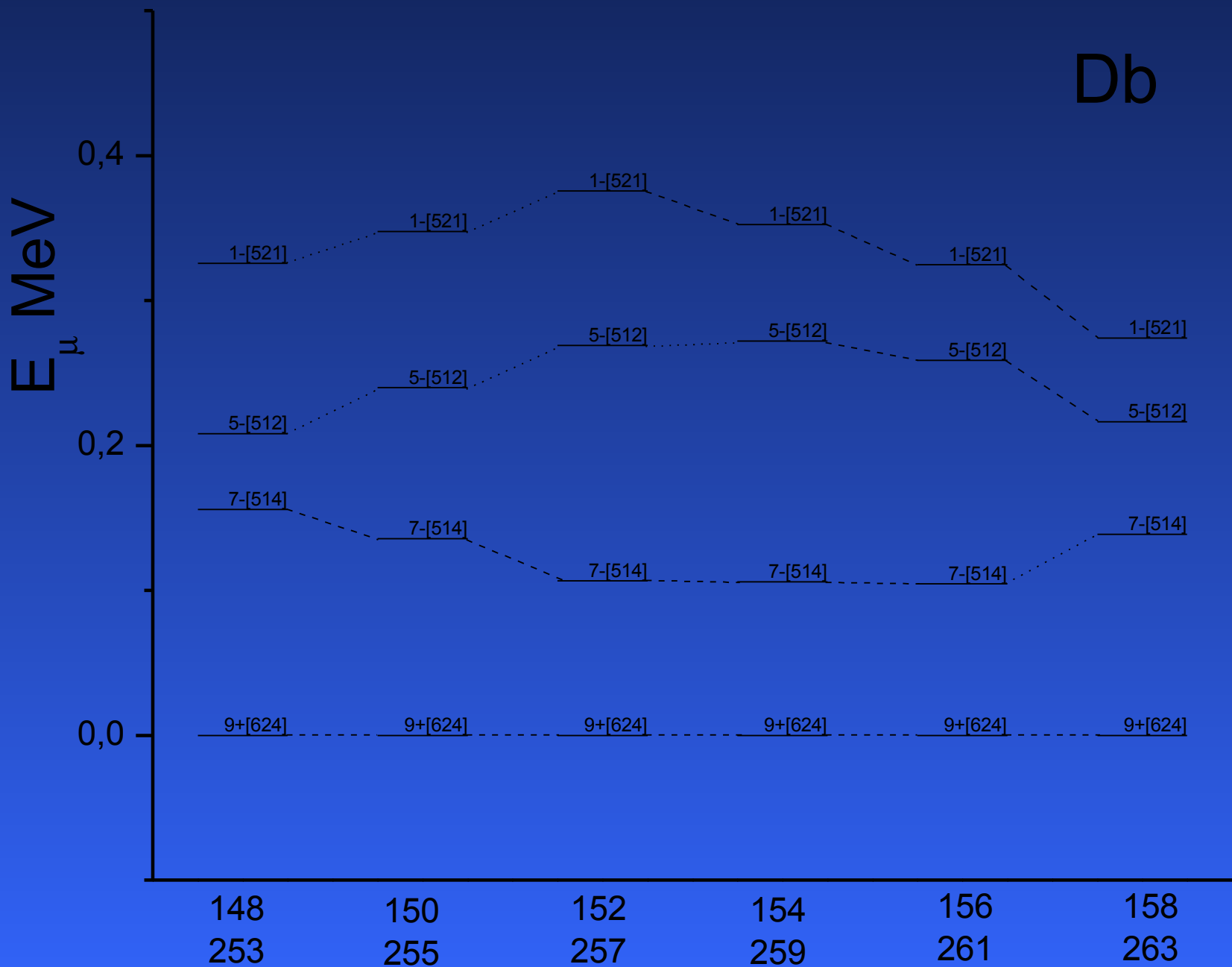








Db

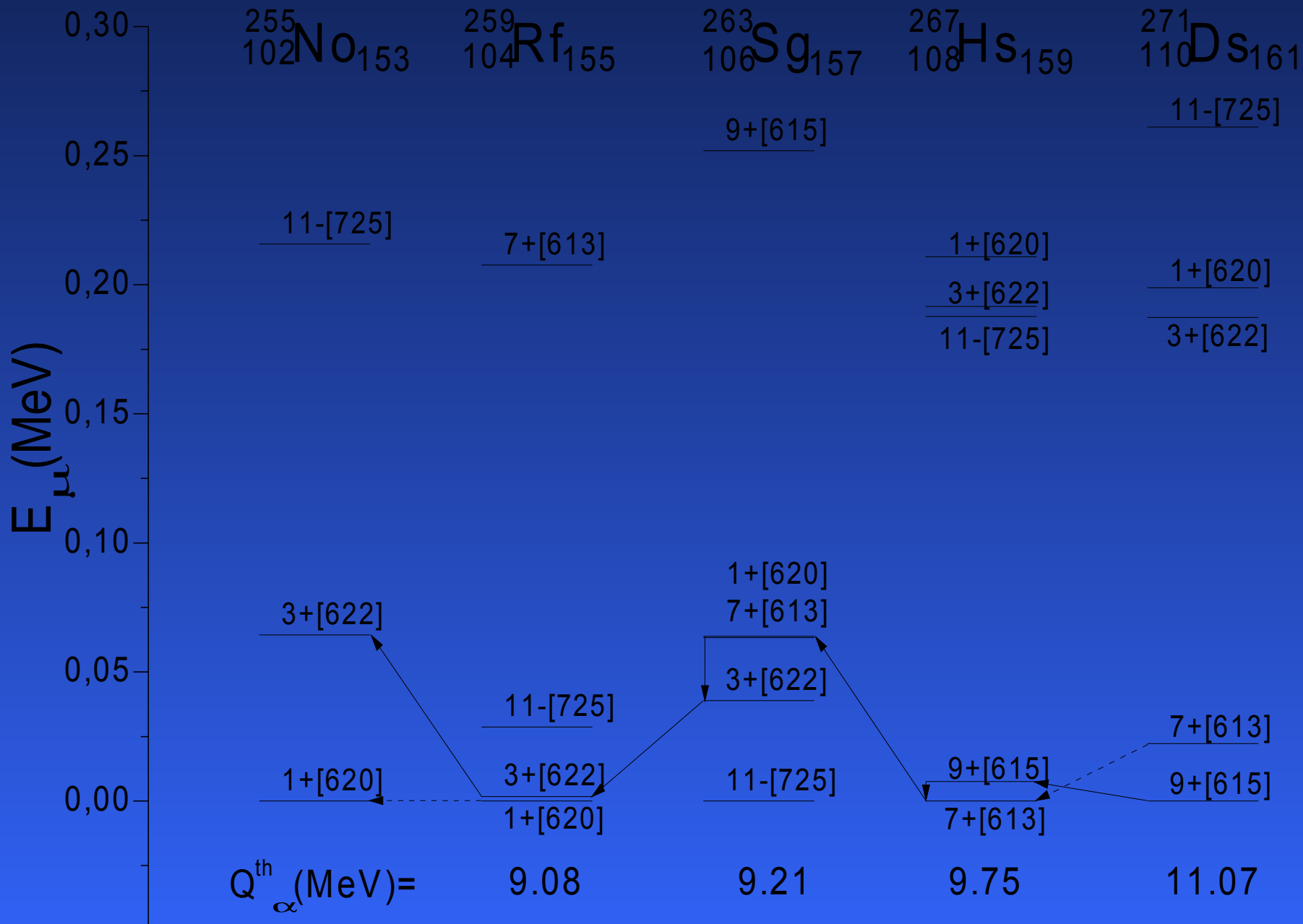


Theoretical and experimental spins in a ground state

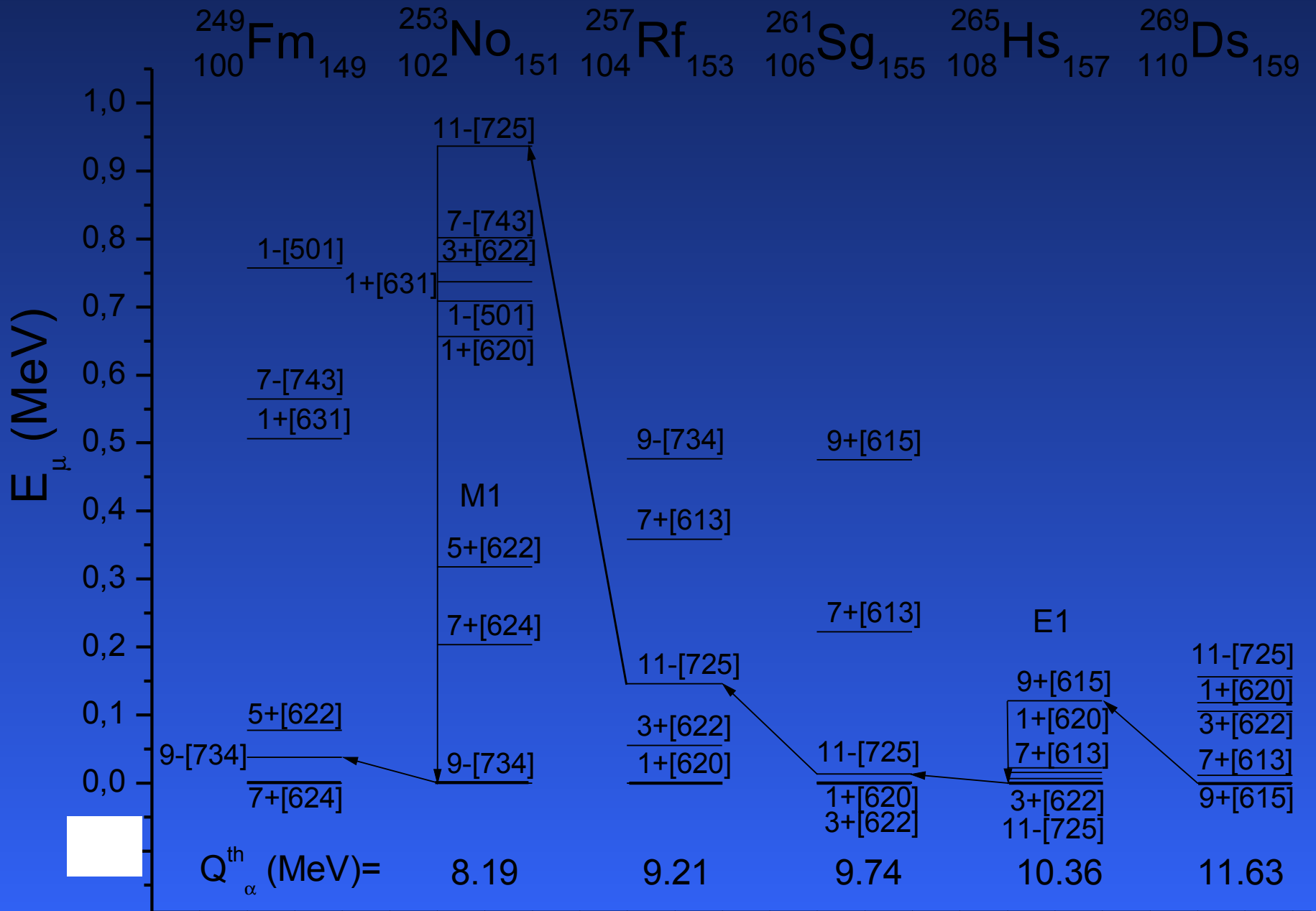
Z=93		
N	s.p. (teor)	s.p. (eksp)
138	5[642]	(5)
140	5[642]	(5+)
142	5[642]	5+
144	5[642]	5+
146	5[642]	5+
148	5[642]	(5+)
150	5[642]	(5-)

Z=95		
N	s.p. (teor)	s.p. (eksp)
142	5[523]	5(-)
144	5[523]	(5)-
146	5[523]	5-
148	5[523]	5-
150	5[642]	(5+)
152	3[521]	(5)

Z=105		
N	s.p. (teor)	s.p. (eksp)
152	9[624]	(9+)



	T_{α}^{exp}	T_{α}^{th}	Q_{α}^{exp} MeV	Q_{α}^{th} MeV	Q_{α}^{trth} MeV
²⁷¹ Ds	1.1 ms	0.39 ms	10.91	11.07	11.06
²⁶⁷ Hs	59 ms	0.25 s	10.03	9.75	9.69
²⁶³ Sg	0.31 s	0.93 s	9.39	9.21	9.25
²⁵⁹ Rf	3.1 s	0.59 s	9.03	9.08	9.08



	T_{α}^{exp}	T_{α}^{th}	Q_{α}^{exp} MeV	Q_{α}^{th} MeV	Q_{α}^{trth} MeV
²⁶⁹ Ds	179 μs	40 μs	11.28	11.63	11.51
²⁶⁵ Hs	701 μs	4.7 ms	10.35	10.68	10.36
²⁶¹ Sg	61 ms	96 ms	9.68	9.74	9.60
²⁵⁷ Rf	9.4 s	64 s	8.80	9.21	8.42
²⁵³ No	92.4 s	90.8 s	8.21	8.19	8.15

Summary

- Quantum characteristics of most experimentally known ground states of analyzed here odd-A nuclei (in 24 of 38 cases for proton-odd and in 23 for 34 cases for neutron-odd nuclei) are reproduced by the calculations.
- Energy of known lowest nucleon excited states are reproduced by the applied model within an average accuracy of about 200 keV.
- Sensitivity of single-particle excitation energies to changes of such a quantity as the equilibrium deformation of a nucleus is rather large.
- A new, simple phenomenological formula for description of α -decay half-lives of heaviest e-e, o-e, e-o and o-o nuclei uses only 5 adjustable parameters. The formula allows one to describe T_{exp} of 61 e-e nuclei roughly within a factor of 1.3, 45 o-e nuclei within a factor of 2.1, 55 e-o nuclei within a factor of 3.2 and 40 o-o nuclei within a factor of 4.0, on the average, when Q_{α}^{exp} is taken.
- The accuracy of the mentioned above phenomenological formula decreases by a factor of about 4, when theoretical values of Q_{α}^{th} are used.
- It is found that description of the ^{289}Ds and ^{271}Ds chains data is rather good.